

METEORS: A DELIVERY MECHANISM OF ORGANIC MATTER TO THE EARLY EARTH

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Abstract. All potential exogenous pre-biotic matter arrived to Earth by ways of our atmosphere, where much material was ablated during a luminous phase called "meteors" in rarefied flows of high (up to 270) Mach number. The recent Leonid showers offered a first glimpse into the elusive physical conditions of the ablation process and atmospheric chemistry associated with high-speed meteors. Molecular emissions were detected that trace a meteor's brilliant light to a 4,300 K warm wake rather than to the meteor's head. A new theoretical approach using the direct simulation by Monte Carlo technique identified the source-region and demonstrated that the ablation process is critical in the heating of the meteor's wake. In the head of the meteor, organic carbon appears to survive flash heating and rapid cooling. The temperatures in the wake of the meteor are just right for dissociation of CO and the formation of more complex organic compounds. The resulting materials could account for the bulk of pre-biotic organic carbon on the early Earth at the time of the origin of life.

Keywords: Ablation, astrobiology, exobiology, meteors, meteoroids, origin of life

1. Introduction

Accretion of extraterrestrial matter has long been of interest as a source of pre-biotic organic carbon for the origin of life on Earth (Oró, 1961; Sagan, 1974; Lewis *et al.*, 1979; Anders, 1989; Pepin, 1991; Huebner and Boice, 1992; Delsemme, 1992; Chyba and Sagan, 1992, 1998; Oberbeck and Aggarwal, 1993; Chang 1993, Whittet, 1997; Oró and Lazcano, 1998). Most organic carbon is thought to be accreted by impacts of comets and primitive asteroids (Oró, 1961; Chyba *et al.*, 1990). However, the high-speed impacts are expected to be so energetic that the ensuing fireball destroys virtually all molecular species in the impacting object and subsequent synthesis upon cooling in a CO₂ rich atmosphere is not efficient (McKay and Borucki, 1997; Chyba and Sagan, 1998). Moreover, such impacts are infrequent and have strong perturbing effects on the pre-biotic environment at the time of delivery (Maher and Stevenson, 1988; McKinnon, 1989; Chyba, 1993).

Accretion of meteoroids is a more gentle and continuous mechanism for delivery of organic carbon species, especially in the case of the Interplanetary Dust Particles (IDP) that are collected mostly intact in the Earth's atmosphere. Among this collected debris of small solar system bodies there are chondritic IDPs with organic carbon abundances of about 10 mass percent on average (Anders, 1989; Gibson, 1992). The organic carbon contains complex aromatic molecules up to 500 a.m.u. (Clemett *et al.* 1993). However, all IDPs represent a small fraction (< 8 %) of the incoming mass, while low encounter velocities favor asteroidal particles that are relatively poor in organic carbon compared to cometary matter (Bradley *et al.*, 1988). Comet Halley's dust has been measured to contain up to 50 mass percent refractory organics (Krueger and Kissel, 1987; Greenberg, 2000).

Here, we consider the possibility that meteors could have delivered organic carbon to the early Earth. Meteors are the luminous phase that represents the process of ablation and fragmentation of the meteoroids' interaction with the atmosphere. Of particular interest are those meteoroids that are too small to cause destructive high-pressure and high-temperature shock-induced chemistry (Menees and Park, 1976, Park and Menees 1978).

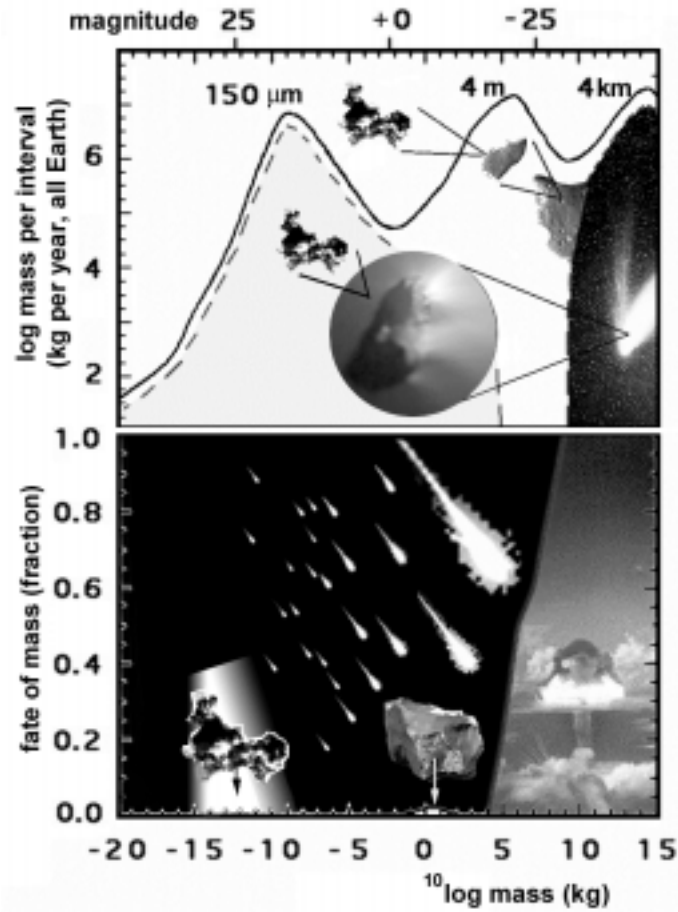


Figure 1. Top panel: the current annual (log) mass influx per unit (log) mass interval of all types of incoming matter in Earth's atmosphere, compiled from Love and Brownlee (1993) and Ceplecha (1992, 1994) Bottom panel: the fate of this matter upon accretion, as a fraction of the total. Fractions are derived from estimated meteorite yields (Obserst *et al.*, 1998; Halliday *et al.*, 1984; Bland *et al.*, 1996), impact limit (Ceplecha, 1992, 1994), and entry heating (Flynn, 1989; Love and Brownlee, 1991), the latter for the new present-day velocity distribution (Taylor and Elford, 1998).

The fraction of mass of meteoroids that are too large and too fast to survive ablation, but too small to cause shock-induced modification of the incoming meteoroid, is schematically shown in Figure 1. The figure shows the present annual mass influx per mass interval. The various

sources of matter are identified. Impacts become progressively infrequent for more massive objects. The present-day distribution is thought to be representative of the dust influx on the early Earth, except that the relative contributions from asteroids and comets may have differed (Chyba and Sagan, 1998). The mass distribution has peaks at 150 microns (meteoroids), 4 meters (asteroid fragments and comet boulders), and 4 km (comets and asteroids). Meteoroids survive as IDPs or ablate during interaction with the atmosphere while in the molecular flow regime, and are recovered as micrometeorites that are mostly CM-type materials and lesser amounts of CI-type materials (Kurat *et al.*, 1994; Engrand and Maurette, 1998). The shaded area reflects organic-rich cometary dust. The meter-sized meteoroids lose most of their mass in the continuum flow regime, where they develop shocks because of their subsonic entry speeds. Some asteroidal fragments will survive as meteorites. Impacting comets and asteroids will catastrophically fragment in an airburst or explode upon impact (Chyba *et al.*, 1990).

Based on this mass influx distribution, and the fate of accreting matter, we submit that meteoroids and small comet boulders can account for the bulk of the organic carbon on the early Earth – that is, if the survival of exogenous organic carbon, or the creation of reduced molecules by atmospheric chemistry, is efficient. Little is known about whether organic carbons can survive the meteor phase. Conditions of free molecular flow (in case the mean free path in air is larger than the typical dimensions of the object) and high Mach number conspire against theoretical and laboratory studies of meteoric plasmas. Factual information relies on remote sensing, to which few modern techniques have been applied.

Here, we report measurements of air plasma temperatures in the wake of meteors, which are compared to models of air plasma emission. The data are interpreted by means of direct simulation by Monte Carlo modeling of meteoroids in the free molecular flow regime. We conclude that meteors offer interesting pathways for the survival of organic carbon that warrant further study.

2. Spectroscopy of Meteoric Plasmas

We deployed from an airborne platform a new un-intensified slit-less CCD spectrograph for near-infrared and visible wavelengths at the time of the intense 1998 Leonid shower during the 1998 Leonid Multi-

Instrument Aircraft Campaign (Jenniskens and Butow, 1999). The spectrograph consisted of a 600 l/mm grating, a Nikon f2.8/300 mm lens, and a Pixelvision two-stage thermoelectrically cooled 1024 x 1024 pixel CCD camera. The same spectrograph was used again to probe Perseid meteor spectra from a ground site in August of 1999. Our best spectra probe 1 cm-sized meteoroids with entry velocities of 61 km/s (Perseids) and 72 km/s (Leonids) at altitudes 90–100 km. The spectra cover part of the wavelength range 580–900 nm at a relatively high 0.5 nm resolution (Figures 2 and 3).

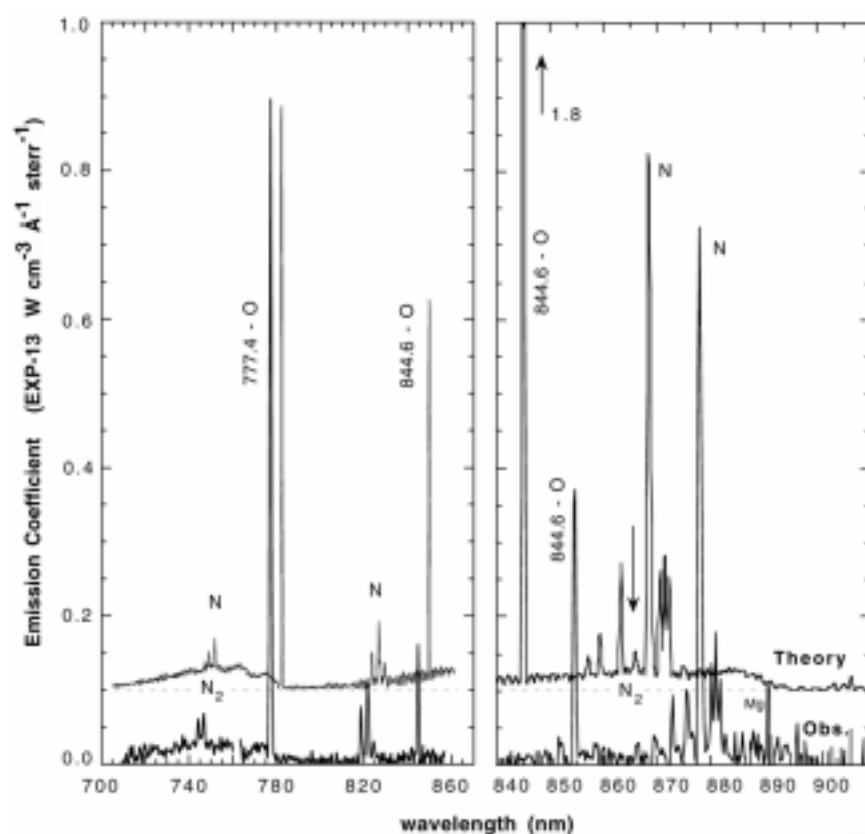


Figure 2. Two Leonid spectra from Nov. 17, 1998, at 17:47:06 UT (left) and 18:08:47 UT (right). The spectra are compared with NEQAIR2 model calculations, which is slightly displaced to facilitate comparison. Note the different line intensities of the OI line emission at 844.6 nm and the NI line emission at 865.6 nm (arrow). The line at 880.7 nm is of MgI.

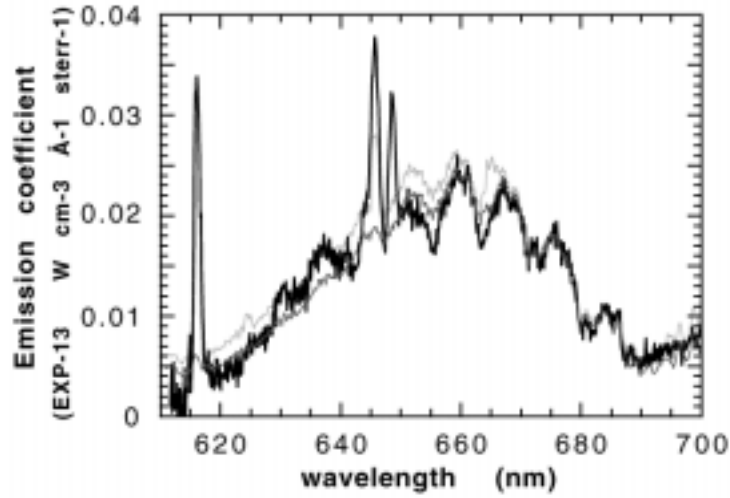


Figure 3. The $\Delta v = 3$ first positive band of N_2 as resolved in the spectrum of a magnitude -1 Perseid meteor from 00:25:26 UT, Aug. 13, 1999 (dark line). Emission lines at 615.7, 626.2, 645.4, and 665.4 nm are from atomic oxygen. The line at 648.9 nm may also be of OI. The two gray lines show the NEQAIR simulations at 4,300 K (dark, down) and 4,670 K (light, up), respectively. The simulation considers pure air (79% N_2 and 21% O_2) at the standard atmosphere pressure of $P = 10^{-6}$ atm at the altitude of the meteor (95 km). The spectrum was convolved with a triangular slit function of FWHM = 0.5 nm, which matches the line width measured on the observed spectra.

The dominant emission is from atmospheric lines of O and N and the first positive bands of N_2 , in contrast to prior studies at shorter wavelengths where the emission lines of ablated meteoric metal atoms dominate. The observed lines and bands are well matched by the NEQAIR2 radiation model of heated air in thermodynamic equilibrium (Park, 1985; Laux, 1993). The match implies that the bulk of emission is from gas in near thermal equilibrium despite the high Mach number flow. The observed ratio of atomic and molecular nitrogen in the Leonid spectra (Figure 2) is a sensitive measure of temperature and implies a chemical equilibrium temperature of $T_c = 4,340 \pm 100$ K. The N_2 band contour (and NI lines) of the Perseids are well matched by a simulation at $T_v = 4,300 \pm 40$ K (Figure 3). All values are similar to temperatures estimated from meteoric metal atom emission lines at $T_e = 4,500 \pm 500$ K (Borovicka *et al.*, 1999).

The data are sufficiently precise to recognize numerous signs of non-equilibrium behavior. There is excess emission at high v levels in the N_2 molecular band as a result of recombination processes (Figure 3). The OI line intensities are not always well matched (Figure 2). Notably, the OI line at 844.6 nm is a factor of 3 fainter than that calculated in all Leonid spectra and different from laboratory LTE air plasmas (Park *et al.*, 1997, 1998), while OI lines between 600 and 700 nm are stronger (Figure 3). Also, the NEQAIR model with initial 0.03% atmospheric CO_2 predicts CN emission comparable to the first positive N_2 bands, but no CN (or isoelectric N_2^+) is observed.

Another sign of non-equilibrium is the intensity of the OI 777.4 nm line emission. The measured intensity of the line, relative to the intensity of recorded background field stars, can be compared to the volume emission coefficient of the plasma in the model (vertical scales in Figure 2). From this, we derive a volume for the emitting gas of $1 \times 10^{13} \text{ cm}^3$, assuming the gas is at 4,300 K. Initial-train-radius theory predicts a volume of only about $3 \times 10^7 \text{ cm}^3$ for the head of the meteor (Jones, 1995).

3. Theoretical Model of rarefied Flow

The source region of the $T \sim 4,300 \text{ K}$ emission was identified using the direct simulation Monte Carlo (DSMC) technique, which was applied to the two-dimensional flow about a 1 cm-sized Leonid meteor (density 1 g/cm^3). For many years DSMC has been developed and applied to a variety of rarefied flows (Dietrich and Boyd, 1996), but this is the first attempt to apply the technique to a computation of this type. Two cases were considered: one with no ablation, and another one with a simple Bronsten ablation model (Bronsten, 1983), in which the ablated material is assumed to be magnesium. In the case without ablation (Figure 4a), we find a rapid decline of temperatures outward, where multiple collisions quickly stop the accelerated air molecules. This is the process described by the initial meteor-train-radius theory, which in light of this model might still apply well to the near spherical source of radar head echoes (Jones, 1995).

Including meteor ablation result in dramatic changes when applying this model (Figure 4b). We find that ablation increases the flow field temperature around the meteor over an extended area in a wake behind the meteor, with elevated values around 5,000 K.

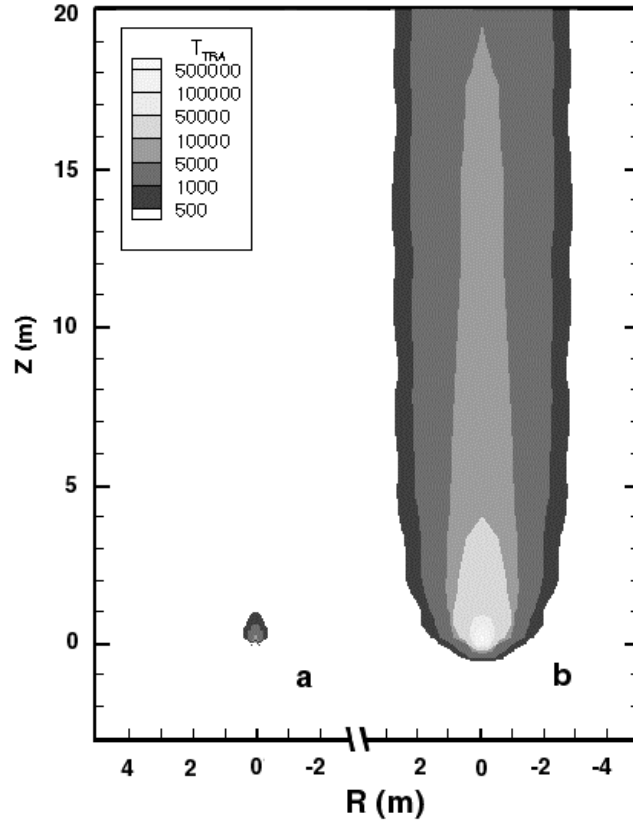


Figure 4. Translational temperature field from a rarefied flow model of a -1 magnitude Leonid meteor at 95 km altitude shown without ablation (a), and with ablation of Mg atoms (b).

The wake is caused by air molecules penetrating the skin of a dense plasma of ablated material in front of the meteoroid that collides with the meteoric plasma several times in a process of thermalisation. The collisions drive the meteoric ablation products and heated air past the meteoroid, where they expand into the meteor wake. The rotational temperatures are typically less than translational temperatures in the meteor's head but they equilibrate in the wake. We conclude that most of the observed meteor emission originates from this nearly equilibrated gas.

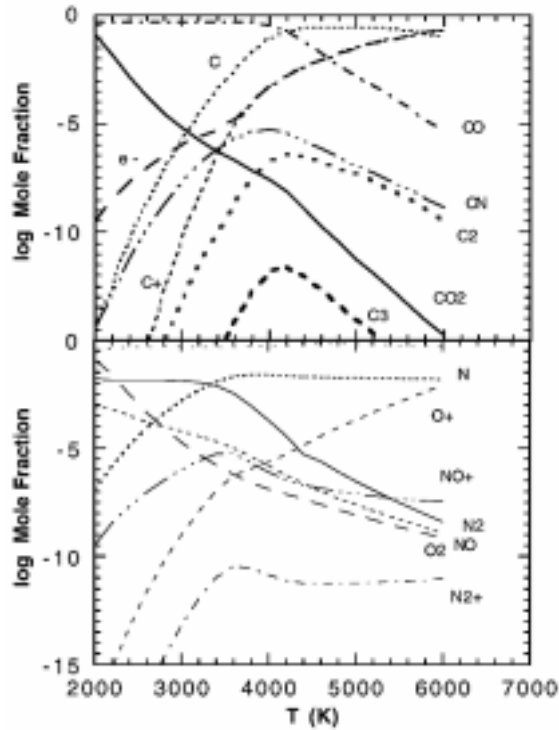


Figure 5. Molecular abundances for equilibrium air plasma at 95 km altitude ($P = 10^{-6}$ atm) in a range of Local Thermodynamic Equilibrium temperatures and for an assumed Mars-like early-Earth atmosphere of particle number composition $O_2/CO_2/N_2/Ar/CO = 0.13/95.32/2.7/1.6/0.08\%$.

4. The Delivery of Organic Matter to the early Earth.

Interestingly, the study of fast shower meteors can help clarify a role of meteors in creating pre-biotic conditions on Earth, which involved a wide range of meteoroid masses and entry velocities. This is because the wake temperatures of all meteors are in the same narrow range of $3,900 \pm 900$ K as derived from the well studied meteoric metal atom emission lines (Borovicka and Bocek, 1995; Borovicka and Betlem; 1997; Harvey, 1973). There is no obvious trend with meteor magnitude (mass) or entry velocity. The observed Leonid spectra, too, do not change significantly with altitude or meteor brightness over the observed range.

We submit that the mixing time-scale of air and ablation plasma does not vary much with meteor mass and velocity. Meteors represent two sources of pre-biotic carbon: a) Direct influx of organic carbon and metallic compounds in a rarefied high Mach number flow, and b) kinetic energy induced atmospheric chemistry involving CO₂ and N₂, which is also of interest to the issue of nitrogen fixation (Chyba and Sagan, 1998).

Our observations show that relevant atmospheric chemistry can occur in two regimes: (1) the extended wake of the meteors at temperatures of about 4,300 K and (2) at the interface layer between impinging air and ablation products at temperatures of about 10,000 K. The Mg⁺ emission line in bright meteors was traced to a component that increases in relative intensity at T~10,000 K with meteor brightness and entry velocity (Borovicka, 1994). This component is due to meteoric vapor that builds up in front of the ablating meteoroid, rather than to the formation of a shock front as is commonly believed. In both cases, the not fully equilibrated chemical reactions occur at higher temperatures and longer time scales than in the 1-D equilibrium models by Menees and Park (1976) and Park and Menees (1978).

Unfortunately, no models are yet capable of reliably handling the types of non-equilibrium chemistry implied by our observations. An important find, however, is that the observed excitation temperatures are close to the dissociation equilibrium of CO. As a result, our equilibrium chemistry simulation of the meteor plasma in a CO₂ rich atmosphere (Figure 5) results in relatively high yields of potential pre-biotic molecules. This air composition may reflect that of the early Earth (Chyba and Sagan, 1998) and is certainly the least favorable case for reaction chemistry. Right at about 4200 K is where the production of linear carbon chains such as C₂ and C₃ peaks. Under these conditions, small amounts of aromatic hydrocarbons are expected to be formed upon cooling, as well as compounds rich in C=O and C-N groups. Such compounds could offer numerous chemical pathways to yield other reduced molecules of potential significance to the origin of life on the early Earth.

Finally, assuming organic compounds are common to Leonid meteoroids, the lack of observed C, C₂ and CN emission from the combustion of organic matter in the ablating meteoroids implies that organic compounds survive as large molecular fragments. It is possible that they are lost early in the meteor trajectory at low temperature T < 500 K when no optical emission occurs. However, meteors do not typically show differential ablation (Borovicka *et al.* 1999). This

suggests that all mineral compounds are lost simultaneously irrespective of volatility in a process rather like sputtering by impinging meteoric ablation products instead of complete evaporation of the grain. Hence, most organic carbon, too, is expected to be lost at 80–110 km altitude under the physical conditions described in this paper.

Perhaps the survival of organic carbon is an analog with the common technique of laser induced desorption of large molecules, whereby they are only momentarily heated in a chemically reducing environment and subsequently quickly cooled by collisions with gas molecules and radiative cooling. For example, laser pulsed heating of aniline on a sapphire surface at 10^8 K/s leads to a peak surface temperature of 1,000 K, but with aniline remaining intact at 360 K internal vibrational temperature (Maechling *et al.*, 1996).

5. Relevance of Delivery through Meteors

At the time of the origin of life on Earth 4 Gyr ago, the mass influx of meteoroids was about 200 times the present day mass influx of 4×10^7 kg/yr (Love and Brownlee, 1993). The factor of 200 follows from the lunar impact record and a linear scaling between meteoroids and parent bodies (Chyba and Sagan, 1998). Approximately half of this influx is thought to be carbon-poor asteroidal matter with at best 4 mass % organic carbon in rare CI chondrites and the other half is carbon-rich (25–50 %) cometary matter (Krueger and Kissel, 1987; Delsemme, 1991). Hence, meteors contributed at least 1×10^9 kg/yr of organic carbon to the early Earth if all organic carbon survived the ablation process.

This outweighs the yield from all other exogenous and terrestrial sources of organic carbon on the early Earth as estimated in the recent review by Chyba and Sagan (1998). The main exogenous source, interplanetary dust particles, was estimated to yield at best only 2×10^8 kg/yr, because less than 8 mass percent of small <1 gram particles survive atmospheric entry heating (Figure 1). Meteorites are a negligible source because only 0.04 mass percent of matter >1 g arrives to Earth in the form of meteorites (Oberst *et al.*, 1998; Halliday *et al.*, 1984; Bland *et al.*, 1996). For even larger bodies, only negligible organic carbon survives the shock chemistry (Chyba and Sagan, 1998). Hence, meteors could account for at least 80% of all exogenous organic carbon.

A relatively high production yield of C_2 and CN from atmospheric chemistry is derived from Figure 4. For a CO_2 rich atmosphere, we

calculate a yield of carbon of 1×10^{-9} kg/J. For a total kinetic energy of 1×10^{18} J/yr from meteors 4 Gyr ago, the yield is 1×10^9 kg/yr. This is only an upper limit because the chemical reactions are not expected to reach full thermodynamic equilibrium, which is supported by the absence of CN or C₂ emission in the observed spectra.

On the other hand, alternative scenarios imply low yields as well. From organic residues at the K/T boundary, the yield of organic carbon in giant impacts 4 Gyr ago in a mildly reducing atmosphere has been optimistically estimated at only 2×10^8 kg/yr (Chyba and Sagan, 1998).

Terrestrial sources are not much more efficient in this regard. Electric discharges contribute a mere 3×10^7 kg/yr and UV photolysis by sunlight in a mildly reducing atmosphere can account for 3×10^8 kg/yr (Chyba and Sagan, 1998). Furthermore, the analysis of effluents from biologically uncontaminated hot or cold springs has failed to show any organic molecules besides methane (Mojzsis *et al.*, 1999).

That leaves meteors as a potential source for more than 2/3 of the prebiotic organic carbon on the early Earth.

Of particular interest for future studies of meteors is the possible detection of ablated organic carbon by its 3.4-micron feature, as molecules in stratospheric aerosols, or indirectly from its interaction with the airglow chemistry. Additional Leonid storms are anticipated in November 2001 and 2002 (McNaught and Asher, 1999), when meteors may fill the sky again as frequently as they did at the time of the origin of life.

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